

# Observation of flat band for terahertz spoof plasmon in metallic kagomé lattice

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We study the dispersion relation of a metamaterial composed of metallic disks and bars arranged to have kagomé symmetry and find that a plasmonic flat band is formed by the topological nature of the kagomé lattice. To confirm the flat-band formation, we fabricate the metamaterial and make transmission measurements in the terahertz regime. Two bands formed by transmission minima that depend on the polarization of the incident terahertz beams are observed. One of the bands corresponds to the flat band, as confirmed by the fact that the resonant frequency is almost independent of the incident angle.

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## I. INTRODUCTION

Kagomé lattices have attracted considerable interest from the aspect of geometric frustration in condensed-matter physics.<sup>1</sup> There is extensive degeneracy of nondispersing resonant modes in a resonator system with kagomé symmetry. These eigenmodes form a flat band, where the resonant frequency of the band is the same for all wavevectors in the first Brillouin zone. The flat band originates purely from the topology of the lattice structure, and it remains flat even when the couplings between the adjacent sites are significantly large. Furthermore, such flat bands can lead to ferromagnetism of itinerant fermions,<sup>2–5</sup> supersolidity for bosons,<sup>6,7</sup> crystalline ordering,<sup>8</sup> and other effects.

Although the flat-band formation is first expected in quantum systems, it can occur in electromagnetic systems. The presence of electromagnetic flat bands in kagomé lattices has already been predicted theoretically in some electromagnetic systems, such as two-dimensional photonic crystals<sup>9</sup> and metallophotonic waveguide networks.<sup>10</sup> In the flat band, the group velocity is slowed down in all directions and the effective mass of the photons becomes very heavy. It is important to study the flat band in the electromagnetic system with kagomé symmetry in terms not only of fundamental physics, but also from an application standpoint, such as slow light;<sup>11</sup> however, there has been no experimental demonstration for the electromagnetic flat band.

Here, we focus on the flat band for a terahertz (THz) plasmonic mode. Although there is no surface plasmon of metals in the THz region, the modes having the dispersion relation similar to surface plasmons are formed in structured metals, and called spoof surface plasmons.<sup>12–15</sup> In this paper, we theoretically and experimentally obtain the dispersion relation for a spoof plasmon in the metallic kagomé lattice and demonstrate the electromagnetic flat band in the THz regime. Numerical simulations are also performed to provide confirmation of the experiments.

## II. THEORETICAL MODEL

We introduce kagomé-type bar-disk resonators (KBDRs), shown in Fig. 1(a). Metallic disks and narrow bars are connected to form a kagomé lattice. KBDRs are artificially engineered metallic structures, and considered as a kind of metamaterial.<sup>16–23</sup> In KBDRs, electric charge stored on each disk temporally oscillates between positive and negative values. We discuss the formation of a flat band in KBDRs by using a coupled oscillator model. We denote the electric potential at the  $i$ th disk as  $\phi_i$ . Introducing  $\Phi_i := \int \phi_i dt$ , we obtain a Lagrangian  $\mathcal{L}$  as

$$\mathcal{L} = \frac{C}{2} \sum_i \dot{\Phi}_i^2 + C_M \sum_{i,j} \frac{A_{ij}}{2} \dot{\Phi}_i \dot{\Phi}_j - \frac{1}{2L} \sum_{i,j} \frac{A_{ij}}{2} (\Phi_i - \Phi_j)^2, \quad (1)$$

with capacitance  $C$  of the disk, inductance  $L$  of the bar, coefficient of electric induction  $C_M$  between nearest disks, and adjacency matrix  $A_{ij}$  of the kagomé lattice, whose element is 1 if the  $i$ th and  $j$ th disks are directly connected by a bar for  $i \neq j$ ; otherwise 0.<sup>24</sup> The first, second, and third terms of Eq. (1) represent the electric energy stored on disks, mutual electric energy stored between disks, and magnetic energy stored around bars, respectively. Here, we consider only the nearest mutual couplings.

Using the Euler-Lagrange equation  $(d/dt)(\partial \mathcal{L} / \partial \dot{\Phi}_i) -$

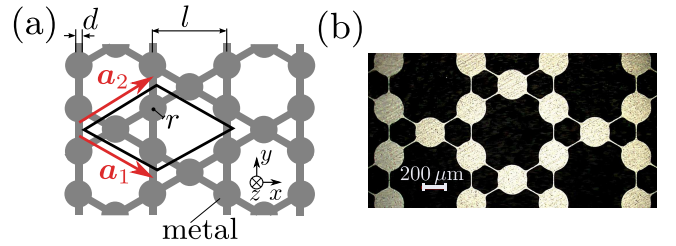


FIG. 1. (Color online) (a) Schematic view of kagomé-type bar-disk resonators (KBDRs). (b) The fabricated KBDRs on a stainless-steel plate with  $l = 800$ ,  $d = 10$ , and  $r = 145$ , and thickness  $h = 30 \mu\text{m}$ .

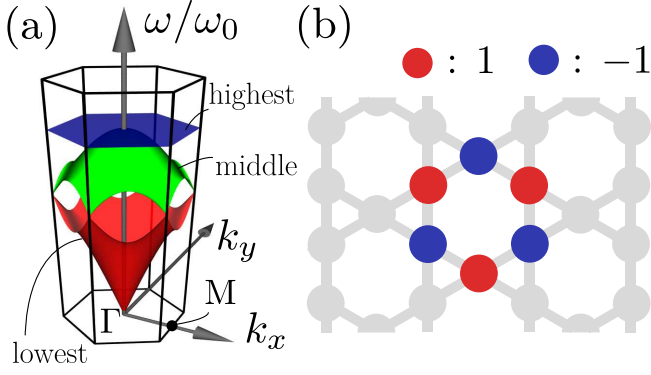


FIG. 2. (Color online) (a) Dispersion relation of KBDRs for  $\kappa = 0$ . (b) Localized eigenmode of KBDRs.

$\partial\mathcal{L}/\partial\Phi_i = 0$ , we obtain coupled charge equations as

$$\ddot{q}_i + \omega_0^2(4q_i - \sum_j A_{ij}q_j) + \kappa \sum_j A_{ij}\ddot{q}_j = 0, \quad (2)$$

with stored charge  $q_i = C\dot{\Phi}_i$  at the  $i$ th disk, resonant angular frequency  $\omega_0 = 1/\sqrt{LC}$ , and coupling coefficient  $\kappa = C_M/C$ . In the frequency domain, we rewrite Eq. (2) as

$$\sum_j A_{ij}\tilde{q}_j = \frac{4 - (\omega/\omega_0)^2}{1 + \kappa(\omega/\omega_0)^2}\tilde{q}_i, \quad (3)$$

where tildes represent complex amplitudes. Owing to the lattice symmetry, we can reduce Eq. (3) to an eigenvalue problem for a  $3 \times 3$  matrix and obtain the dispersion relation consisting of three bands as

$$\frac{\omega}{\omega_0} = \sqrt{\frac{6}{1-2\kappa}}, \sqrt{\frac{3 + 2(3+F)\kappa \pm (1+4\kappa)\sqrt{3+2F}}{1+2\kappa-2(1+F)\kappa^2}}, \quad (4)$$

where  $F = \cos \mathbf{k}_{\parallel} \cdot \mathbf{a}_1 + \cos \mathbf{k}_{\parallel} \cdot \mathbf{a}_2 + \cos \mathbf{k}_{\parallel} \cdot (\mathbf{a}_1 - \mathbf{a}_2)$  with wavevector  $\mathbf{k}_{\parallel}$  in the  $xy$  plane and unit-lattice vectors  $\{\mathbf{a}_1, \mathbf{a}_2\}$  shown in Fig. 1(a). A calculated dispersion relation is shown in Fig. 2(a) for  $\kappa = 0$ . It is clear that the highest band  $\omega/\omega_0 = \sqrt{6/(1-2\kappa)}$  is flat or independent of  $\mathbf{k}_{\parallel}$ . It can be seen that the lowest band shows conical dispersion near the  $\Gamma$  point and that the bending middle band touches the flat band at the  $\Gamma$  point.

The flat band is caused by the interference of spoof plasmon propagating in the kagomé lattice. The adjacency matrix  $A_{ij}$  of the kagomé lattice has eigenmodes localized at hexagonal sites with an eigenvalue of  $-2$ . One of them is shown in Fig. 2(b). The number of the eigenmodes is equal to the number of hexagons in the kagomé lattice. The flat band is formed from these degenerated localized modes as they are not coupled with each other.

### III. EXPERIMENTAL SETUP

We fabricate KBDRs on a stainless-steel plate. The dimensions depicted in Fig. 1(a) are as follows: period between bars  $l = 800 \mu\text{m}$ , bar width  $d = 10 \mu\text{m}$ , disk radius  $r = 145 \mu\text{m}$ , and metal thickness  $h = 30 \mu\text{m}$ . The size of the area patterned KBDRs is  $1.1 \text{ cm} \times 1.1 \text{ cm}$ . A photomicrograph of a fabricated sample is shown in Fig. 1(b).

To investigate the dispersion relation experimentally, we perform THz time-domain spectroscopy (THz-TDS), shown in Fig. 3. A THz emitter and detector (EKSPLA Ltd.) with dipole antennas attached to Si lenses are used. These antennas are integrated on low-temperature-grown GaAs photoconductors, and driven by a femtosecond fiber laser (F-100, IMRA America, Inc.) with a wavelength of  $810 \text{ nm}$  and pulse duration of  $120 \text{ fs}$ . The THz beam is collimated with the Si lens near the emitter. The beam radius is about  $3.7 \text{ mm}$ , which covers a large number of KBDRs. The THz electric field  $E(t)$  is coherently measured by the detector. We obtain the transmission spectrum  $T(\omega)$  in the frequency domain from  $T(\omega) = |\tilde{E}_s(\omega)/\tilde{E}_{\text{ref}}(\omega)|^2$ , where  $\tilde{E}_s(\omega)$  and  $\tilde{E}_{\text{ref}}(\omega)$  are Fourier transformed electric fields with and without the sample, respectively.

In order to obtain the band structure between the  $\Gamma$  point and the M point, the sample is rotated by  $\theta$  with respect to the  $y$  axis from normal incidence (Fig. 3). The angles  $\theta$  range from  $\theta = 0^\circ$  to  $\theta = 65^\circ$  with a step  $\Delta\theta = 2.5^\circ$ . The magnitude of the wavevector  $\mathbf{k}_{\parallel}$  on the sample plane is given by  $k_{\parallel} = (\omega/c) \sin \theta$ , where  $c$  is the speed of light.

We perform transmission experiments for two polarizations along the  $x'$  axis (parallel configuration) and  $y$  axis (perpendicular configuration). We denote the electric field of the incident wave as  $\mathbf{E}$ , and the projection of  $\mathbf{E}$  to the sample plane as  $\mathbf{E}_{\parallel}$ . For parallel or perpendicular configurations,  $\mathbf{E}_{\parallel}$  is parallel or perpendicular, respectively, to  $\mathbf{k}_{\parallel}$ . Wire-grid polarizers near the emitter and detector are adjusted so that the emitted and detected fields have the same polarization.

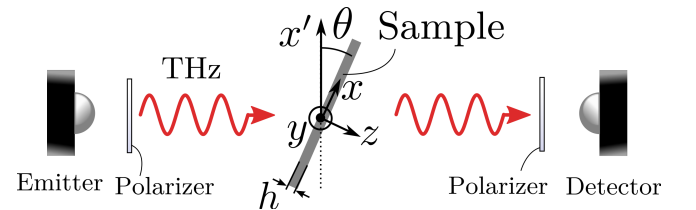


FIG. 3. (Color online) Schematic view of the transmission experiment. The plane of a sample is represented by the coordinate system  $(x, y)$  shown in Fig. 1(a).

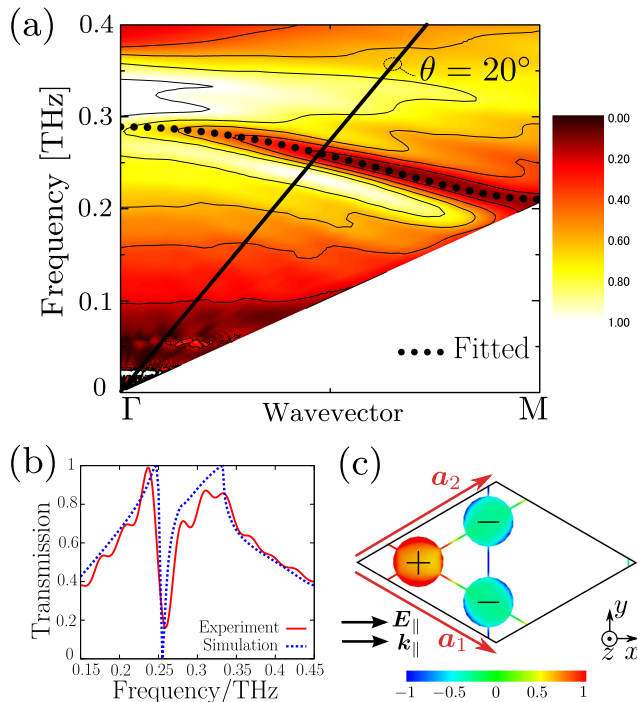


FIG. 4. (Color online) Parallel configuration ( $\mathbf{E}_{\parallel} \parallel \mathbf{k}_{\parallel}$ ). (a) Experimentally obtained transmission diagram of KBDRs. Transmission minima between 0.21 THz and 0.28 THz are observed and theoretically fitted by a dotted line. The solid line corresponds to  $\theta = 20^\circ$ . (b) Transmission spectrum for  $\theta = 20^\circ$  obtained by simulation and experiment. (c) Surface electric charge distribution obtained by simulation at the transmission minimum 0.255 THz for  $\theta = 20^\circ$  in the unit cell.

#### IV. RESULTS

Figure 4(a) displays the transmission spectrum for parallel configuration. The wavevectors are estimated as  $k_{\parallel} = (\omega/c) \sin \theta$ . Transmission spectrum minima are observed from 0.21 THz to 0.28 THz. With an increase of wavenumber, the frequency of the transmission minimum decreases from 0.28 THz at the  $\Gamma$  point and approaches 0.21 THz at the M point.

For further investigation, we calculate the electromagnetic response of KBDRs for parallel configuration. A commercial finite-element method solver (Ansoft HFSS) is used. In the simulation, a plane THz wave is injected into perfectly conducting KBDRs at an incident angle  $\theta$ . By using periodic boundaries with some phase shifts, the transmission and the electromagnetic fields in the unit cell are calculated for an oblique incident plane wave. The measured transmission spectra for  $\theta = 20^\circ$  are compared with the simulation in Fig. 4(b). The frequency of the transmission minimum and shape of the curve for the simulation are consistent with the experimental result, which confirms the validity of the assumption of perfect conductors.

Figure 4(c) shows the calculated distribution of surface electric charges at a minimum (0.255 THz) for  $\theta = 20^\circ$ .

This mode corresponds to the middle band. Disks along the  $x$  axis are alternately charged. The in-phase currents along  $\mathbf{a}_1$  and  $\mathbf{a}_2$  are excited by the electromotive force due to  $\mathbf{E}$ . No resonance appears for  $\theta = 0$  because the current flowing into a disk is balanced by the current flowing out of it. For the excitation of this mode, a phase-shifted field in the  $x$ -direction is needed.

By using Eq. (4), the fitting parameters are obtained from experimental data as  $\omega_0/(2\pi) = 0.105$  THz and  $\kappa = 0.103$ . The resultant dispersion curve is represented as a dotted curve in Fig. 4(a). Positive charges on one disk induce negative charges on the other; therefore,  $\kappa < 0$  is ordinarily expected in the static limit ( $\omega \rightarrow 0$ ). It seems strange that  $\kappa$  would be positive. In our situation, it can be explained by a retardation effect.<sup>25</sup> The phase shift between nearest disks is given by  $(\omega_c/c) \times l/\sqrt{3} \sim 0.77 \times \pi$  at the center frequency  $\omega_c/(2\pi) = 0.25$  THz of the middle band. The near  $\pi$  shift leads to  $\kappa > 0$ . Although  $\kappa$  depends on frequency, we can approximately regard it as a constant between 0.2 THz and 0.3 THz.

Figure 5(a) displays the transmission spectrum for perpendicular configuration. Unlike in the case of parallel configuration, the flat band of the transmission minima is observed around 0.28 THz.

In order to confirm that the flat band is due to the interference of a spoof plasmon, we perform a simulation for perpendicular configuration. A calculated transmission spectrum by simulation for  $\theta = 20^\circ$  is shown in Fig. 5(b) with the experimental data. We see a good agreement in the frequency of the transmission minimum and the shape of the curve. The calculated distribution of surface electric charges at a minimum (0.278 THz) for  $\theta = 20^\circ$  is shown in Fig. 5(c). The resonant mode has antisymmetric amplitudes on the right two disks and there is no charge stored on the left disk. This mode can be constructed by the localized modes shown in Fig. 2(b). Therefore, the flat transmission minima are caused by the topological nature of the kagomé lattice. The mode is excited by anti-phase electromotive force caused by  $\mathbf{E}$  along bars parallel to  $\mathbf{a}_1$  and  $\mathbf{a}_2$ . The electromotive force along vertical bars does not contribute to the storage of charges on disks because the currents flowing into and out of a disk are balanced. In the case of  $\theta = 0$ , the current flowing into a disk is equal to the current flowing out of it and the flat-band mode cannot be excited.

A dotted line in Fig. 5(a) represents the highest band given by Eq. (4) with the previously derived parameters  $\omega_0/(2\pi) = 0.105$  THz and  $\kappa = 0.103$ . It fits well with the minima experimentally obtained. The bend of the flat band caused by coupling to second (or higher) nearest sites is negligibly small, so the assumption of only nearest mutual disk coupling is appropriate.

#### V. DISCUSSION

Our bar-disk resonators (BDRs) are obtained by inverting the metallic area and empty space of the slit-hole

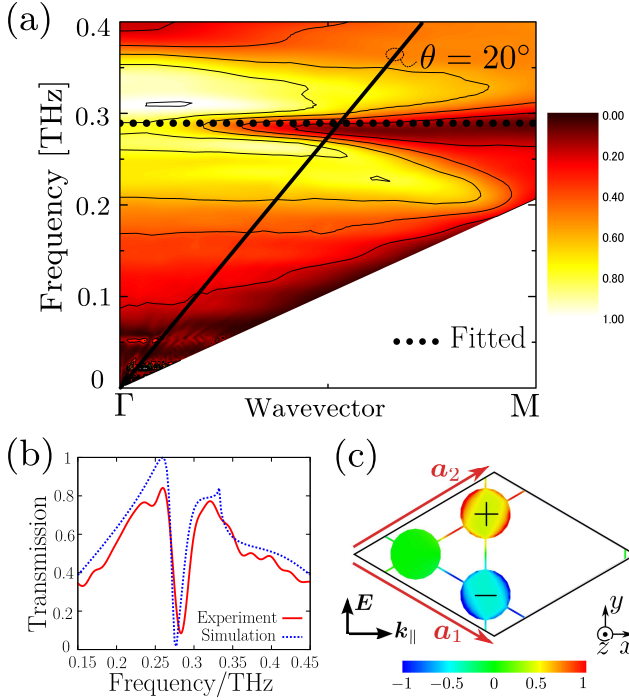


FIG. 5. (Color online) Perpendicular configuration ( $\mathbf{E} \perp \mathbf{k}_{\parallel}$ ). (a) Experimentally obtained transmission diagram of KBDRs. A flat transmission band is observed around 0.28 THz and theoretically fitted by the dotted line. The solid line corresponds to  $\theta = 20^\circ$ . (b) Transmission spectrum for  $\theta = 20^\circ$  obtained by simulation and experiment. (c) Surface electric charge distribution obtained by simulation at a transmission minimum of 0.278 THz for  $\theta = 20^\circ$  in the unit cell.

resonators (SHRs)<sup>26,27</sup> composed of slits and holes engraved on an ultra thin metallic plate. The BDRs and SHRs are complementary structures related through the Babinet's principle,<sup>28–30</sup> based on the electric/magnetic reciprocity of a vacuum. We denote a pair of an electric field  $\mathbf{E}$  and a magnetic field  $\mathbf{H}$  as  $(\mathbf{E}, \mathbf{H})$ . Due to the Babinet principle, the transmittance of a complementary metallic screen with an complementary incident wave  $(\mathbf{E}', \mathbf{H}') = (Z_0 \mathbf{H}, -\mathbf{E}/Z_0)$  is equal to the reflectance of the original metallic screen illuminated by an incident wave  $(\mathbf{E}, \mathbf{H})$ , where  $Z_0$  is the impedance of a vacuum. Thus, the transmission peaks in SHRs correspond to the transmission minima in BDRs. This fact shows the duality of the Lagrangians of SHRs and BDRs.

Electromagnetic flat bands for all crystal directions have been reported for photonic crystals with square symmetry, theoretically<sup>31</sup> and experimentally.<sup>32</sup> In this case, the flat bands are formed due to good lateral confinement (high  $Q$  factor) of the quadrupole modes, which lack preferential coupling directions, at defects of photonic crystals. On the other hand, the flat band for KBDRs is not caused by highly confined modes, but by the topological nature of the kagomé lattice. The

kagomé lattice prevents spoof plasmons from propagating despite the existence of strong coupling. Thus, the physical origin of the flat band on KBDRs differs from that for photonic crystals.<sup>31,32</sup> The flat band for propagating modes has been theoretically predicted for square waveguide networks.<sup>33</sup> Our system is considered as an experimental realization of the flat band for the propagating mode.

The flat band in the kagomé lattice comes from local interference effects. The global symmetry (i.e., periodicity of the lattice) is not necessarily required because local symmetries can support the localized mode. The resonance independent of the incident angle could be expected for the metallic structure having localized modes with the same resonant frequency, even without periodicity.

## VI. CONCLUSION

In conclusion, we studied theoretically and experimentally the electromagnetic flat band on a metallic kagomé lattice. Kagomé-type bar-disk resonators were proposed to realize the flat band. A dispersion relation composed of three bands was theoretically predicted for KBDRs. The highest band was flat for all wavevectors. Two bands formed by transmission minima depending on the polarization of the incident terahertz beams were observed experimentally. One of the bands corresponded to the flat band. Theoretical fitting showed good agreement for these modes. By simulation, we revealed that the flat band was caused by the topological nature of the kagomé lattice.

The flat band can be applicable to slow light, which is useful for the control of group velocity,<sup>34,35</sup> high sensitive sensing, and other applications. In the flat band, the effective mass of photons becomes very heavy and their correlation has an important role. Multiphoton correlation effects in kagomé lattices are important in terms of fundamental physics and should be studied in the future.

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